

## Lecture Outline for Monday, Nov. 4, 2024

## 1. Interpretation of wave equation solution

- a. *Matlab* simulation
- b. Vibration modes and resonances; harmonically related
- c. Standing waves vs. traveling waves. For example, consider the  $g(x) = 0$  (or  $B_n = 0$ ) case. Each eigenmode, which is a *standing* wave, can be expressed as the sum of two counterpropagating *traveling* waves with speed  $v$ :

$$\sin\left(\frac{n\pi x}{L}\right)\cos\left(\frac{n\pi vt}{L}\right) = \frac{1}{2}\sin\left[\frac{n\pi}{L}(x+vt)\right] + \frac{1}{2}\sin\left[\frac{n\pi}{L}(x-vt)\right]$$

All eigenmodes propagate at the same speed  $a$  in this simple example. In more realistic problems, different eigenmodes travel at different speeds (dispersion).

- d. Not modeled in this example:
    - i. Acoustic coupling between strings and nearby non-fixed objects
    - ii. Energy dissipation within string and nearby objects and due to air resistance
    - iii. Scattering (reflections) from nearby objects
    - iv. Presence of nearby objects, especially resonant cavities, which can affect sound perceived by listeners and can alter resonant frequencies somewhat
    - v. Gravity
2. Wave equation (1-D) problems with open boundaries (or boundaries so far away that they can be considered to be infinite in extent)

$$v^2 \frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 u}{\partial t^2} \quad -\infty \leq x \leq \infty, \quad t \geq 0$$

$$\text{ICs: } u(x, 0) = f(x) \quad \text{and} \quad \left. \frac{\partial u(x, t)}{\partial t} \right|_{t=0} = 0$$

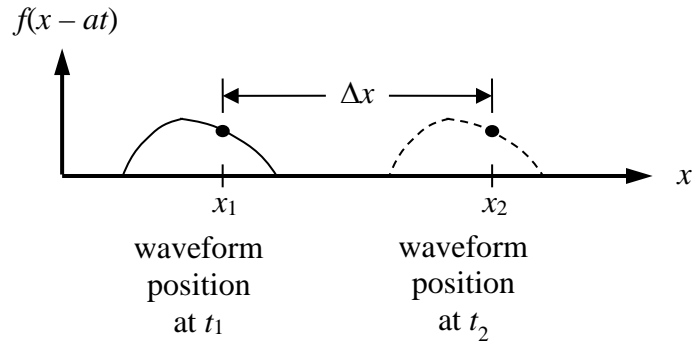
- a. Special solution methods required (e.g., Green's functions for nonhomogeneous problems)
- b. One approach: See supplemental reading on D'Alembert's solution
- c. Solution has  $x \pm vt$  in arguments of functions (traveling waves)

$$u(x, t) = \frac{1}{2}f(x+at) + \frac{1}{2}f(x-at) + \frac{1}{2a}g_A(x+at) - \frac{1}{2a}g_A(x-at),$$

where  $g_A(x)$  is the antiderivative (integral) of  $g(x)$ .

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- d. Example: depiction of propagation of the waveform  $f(x - at)$  over time:



The point on the waveform marked with a dot is at  $x = x_1$  at time  $t_1$ . At later time  $t_2$ , the waveform advances to the right so that the dot reaches the location  $x = x_2$ . The waveform travels the distance  $\Delta x = x_2 - x_1$  over the time interval  $\Delta t = t_2 - t_1$ . The dot represents a particular point along the function  $f(x)$ , so

$$x_2 - at_2 = x_1 - at_1 \quad \rightarrow \quad x_2 - x_1 = a(t_2 - t_1) \quad \rightarrow \quad \frac{x_2 - x_1}{t_2 - t_1} = \frac{\Delta x}{\Delta t} = a \quad (a = \text{speed})$$

3. Wave equation in polar (cylindrical) coordinate system. Example: Vibrating circular membrane of radius  $c$ , assuming only radial vibrational modes (no angle-dependent modes):

$$a^2 \left( \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} \right) = \frac{\partial^2 u}{\partial t^2} \quad \text{for } 0 \leq r \leq c \quad \text{and} \quad t \geq 0$$

- a. ODEs after separation ( $R$  equation is a parametric Bessel equation of order 0)

$$r^2 R'' + rR' + \lambda r^2 R = 0 \quad \text{and} \quad T'' + \lambda a^2 T = 0$$

- b. Boundary conditions and initial conditions and their interpretation

$$u(c, t) = 0, \quad u(r, 0) = f(r), \quad \text{and} \quad \left. \frac{\partial u}{\partial t} \right|_{t=0} = g(r); \quad u \text{ is finite everywhere}$$

- a. Suggested solution forms for ODEs (with  $\lambda_n = \alpha_n^2$ )

$$R_n(r) = c_1 J_0(\alpha_n r) + c_2 Y_0(\alpha_n r) \quad \text{and} \quad T_n(t) = c_3 \cos(a\alpha_n t) + c_4 \sin(a\alpha_n t)$$

- b. The  $R(r)$  problem is an eigenvalue problem (it's also a BVP) and specifically a singular Sturm-Liouville problem; the  $T(t)$  problem is an IVP  
 c. Special additional condition: Solution must be finite within boundary (for  $r \leq c$ )  
 d. In the case of a drum being struck by a stick or mallet,  $f(r) = 0$  and  $g(r)$  is a pulse centered at  $r = 0$ .