Lecture Outline for Monday, Oct. 28, 2024

1. Interpretation of solution

$$
u(x,t) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right) e^{-kn^2\pi^2t/L^2}
$$

- a. How do we break it down?
- b. Does it satisfy BCs?
- c. What can we learn from it?
- d. Does it make sense? Behavior as $n \to \infty$ and $t \to \infty$
- e. Behavior of coefficients in infinite series as $n \to \infty$:

$$
A_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx
$$

- i. As $n \to \infty$, the eigenfunction $\sin(n\pi x/L)$ consists of a large number of cycles that span $x = 0$ to $x = L$.
- ii. If $f(x)$ is relatively smooth, then it is almost constant over each individual cycle of sin(*n*π*x*/*L*).
- iii. The integral from 0 to *L* can be decomposed into a sum of *n* different integrals over the individual cycles of sin(*n*π*x*/*L*).
- iv. Each of those integrals evaluates to almost zero because there are almost equal areas under the curve above and below the *x*-axis. Thus, $A_n \to 0$ as $n \to \infty$.
- f. *Matlab* simulation
- 2. Not always possible to find a solution with the SoV method; some PDE solutions are not separable. Which of the following PDEs can be solved via SOV, and which cannot?

a.
$$
x \frac{\partial u}{\partial x} = y \frac{\partial u}{\partial y}
$$

\nb. $a^2 \frac{\partial^2 u}{\partial x^2} - g = \frac{\partial^2 u}{\partial t^2}$, where *g* is a constant
\nc. $\frac{\partial u}{\partial x} = \frac{\partial u}{\partial y} + u$
\nd. $a^2 \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) = \frac{\partial u}{\partial t}$

(*continued on next page*)

3. Wave equation example

$$
v^2 \frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 u}{\partial t^2} \qquad 0 \le x \le L \,, \quad t \ge 0
$$

- a. Vibrations of string: $v = \sqrt{\frac{T}{\rho}}$, where *T* = tension in string, ρ = mass per unit length
- b. Typical BCs and ICs (two ICs because time problem is second order)

$$
u(0,t) = 0
$$
, $u(L,t) = 0$, $u(x,0) = f(x)$, and $\frac{\partial u}{\partial t}\Big|_{t=0} = g(x)$

c. Apply SOV method with $u(x, t) = X(x)T(t)$ to obtain

$$
v^2 X''T = XT'' \quad \to \quad \frac{X''}{X} = \frac{T''}{v^2 T} = -\lambda \quad \to \quad X'' + \lambda X = 0 \quad \text{and} \quad T'' + v^2 \lambda T = 0
$$

d. The BCs $u(0, t) = 0 \rightarrow X(x) = 0$ and $u(L, t) = 0 \rightarrow X(L) = 0$ lead to the orthogonal spatial eigenfunctions

$$
X_n(x) = c_2 \sin\left(\frac{n\pi x}{L}\right) \quad \text{with} \quad \lambda_n = \left(\frac{n\pi}{L}\right)^2 \quad \text{for} \quad n = 1, 2, 3, \dots
$$

e. The *T* problem is an initial value problem involving the Fourier equation. The possible solution forms are

$$
T_n(t) = c_4 \cos\left(\nu \sqrt{\lambda_n} t\right) + c_5 \sin\left(\nu \sqrt{\lambda_n} t\right) = c_4 \cos\left(\frac{n\pi vt}{L}\right) + c_5 \sin\left(\frac{n\pi vt}{L}\right)
$$

or
$$
T_n(t) = c_4 e^{iv\sqrt{\lambda_n} t} + c_5 e^{-iv\sqrt{\lambda_n} t},
$$

but the one that uses trigonometric functions is more convenient since it involves only real quantities. Experience shows that this is true even though $t \rightarrow \infty$ (an open "boundary").

f. Each eigensolution to the full PDE with $A_n = c_2c_4$ and $B_n = c_2c_5$ has the form:

$$
u_n(x,t) = X_n(x)T_n(t) = \left[A_n \cos\left(\frac{n\pi vt}{L}\right) + B_n \sin\left(\frac{n\pi vt}{L}\right)\right] \sin\left(\frac{n\pi x}{L}\right)
$$

g. Full solution to the PDE for the general case:

$$
u(x,t) = \sum_{n=1}^{\infty} u_n(x,t) = \sum_{n=1}^{\infty} \left[A_n \cos\left(\frac{n\pi vt}{L}\right) + B_n \sin\left(\frac{n\pi vt}{L}\right) \right] \sin\left(\frac{n\pi x}{L}\right)
$$